On the Bernstein inequality

Samira Alamin Sulaiman^(*)

Dept. of mathematics, Faculty of sciences, University of Zawia

Abstract

Chemin [1], proved the inequality of Bernstein for any tempered distribution u. In this paper, we will extend its proof for a bloc dyadic $\dot{\Delta}_q u$ and $S_q u$. We will use the Fourier transform and apply the Yong inequality for convolution. In addition, we will use the techniques of analysis in frequency space.

Keywords: Dyadic decomposition, Littlewood-Paley operators, radial functions, space of Schwartz, Bernstein inequality.

83

^(*) Email: samira.sulaiman@zu.edu.ly

1. Introduction

In this section, we recall the Young inequality, we define the dyadic decomposition of the full space \mathbb{R}^d and recall the Littlewood-Paley operators.

The following inequality is well-known, can be found for example in [3].

Lemma 1.1: (Young convolution inequality)

For any two functions f and g, such that $f \in L^c$ and $g \in L^a$ and any constants $(a, b, c) \in [1, \infty]^3$, such that

$$1 + \frac{1}{b} = \frac{1}{c} + \frac{1}{a}$$

Then we have $f * g \in L^b$ and

 $||f * g||_{L^b} \le C ||f||_{L^c} ||g||_{L^a}$, *C* is a constant.

We can conclude immediate the following result, we refer to [1], [2].

Lemma 1.2 :

For every function $f \in S$, where S is the space of Schwartz such that $f \in L^1 \cap L^\infty$ and for every $1 < c < \infty$, then we have $f \in L^c$ and $(1 + |\cdot|^2)^d \partial^\alpha f$ is bounded.

We recall the Littlewood-Paley operators see [1], [4] and [6] for more details.

Definition 1.3:

There exist two non-negative radial functions $\chi \in \mathcal{D}(\mathbb{R}^d)$ and $\varphi \in \mathcal{D}(\mathbb{R}^d/\{0\})$ such that

 $\chi(\xi) + \sum_{q \ge 0} \varphi(2^{-q}\xi) = 1, \quad \forall \xi \in \mathbb{R}^d,$

University Bulletin – ISSUE No.22- Vol. (2) – June- 2020.

84

$$\sum_{q \in \mathbb{Z}} \varphi(2^{-q}\xi) = 1, \qquad \forall \xi \in \mathbb{R}^d / \{0\},$$
$$|p - q| \ge 2 \Rightarrow supp \ \varphi(2^{-p}.) \cap supp \ \varphi(2^{-q}.) = \varphi,$$
$$q \ge 1 \Rightarrow supp \ \chi \cap upp \ \varphi(2^{-q}.) = \varphi.$$

Let $h = \mathcal{F}^{-1}\varphi$ and $\bar{h} = \mathcal{F}^{-1}\chi$, the frequency localization inhomogeneous operators Δ_q and S_q are defined by

$$\begin{aligned} \Delta_q f &= \varphi(2^{-q}D)f, \qquad S_q f = \chi(2^{-q}D)f\\ \Delta_{-1} f &= S_0 f, \qquad \Delta_q f = 0 \quad for \ q \leq -2. \end{aligned}$$

And the frequency localization homogeneous operators $\dot{\Delta}_q$ and \dot{S}_q are defined by

$$\dot{\Delta}_q f = \varphi(2^{-q}D)f, \qquad \dot{S}_q f = \chi(2^{-q}D)f$$

We notice that $\Delta_q = \dot{\Delta}_q, \forall q \in \mathbb{N}$ and S_q coincides with \dot{S}_q on tempered distributions modulo polynomials.

From the definition of the operator Δ_q , we can write ([3], [5]),

$$u = \sum_{q} \Delta_{q} u$$

2. Bernstein inequality

In this section, we will prove a Bernstein inequality for a tempered distribution u with a bloc dyadic $\dot{\Delta}_q$ and S_q which is the main result of this paper.

Lemma 2.1: (Bernstein Lemma) There exists a constant C > 0 such that for every $q \in \mathbb{Z}, k \in \mathbb{N}$ and for every tempered distribution u we have

$$sup_{|\alpha|=k} \|\partial^{\alpha} S_{q}u\|_{L^{b}} \leq C^{k} 2^{q\left(k+d\left(\frac{1}{a}-\frac{1}{b}\right)\right)} \|S_{q}u\|_{L^{a}} \quad , b \geq a \geq 1 \dots \dots (1)$$

85

$$C^{-k}2^{qk} \|\dot{\Delta}_{q}u\|_{L^{a}} \leq sup_{|\alpha|=k} \|\partial^{\alpha}\dot{\Delta}_{q}u\|_{L^{a}} \leq C^{k}2^{qk} \|\dot{\Delta}_{q}u\|_{L^{a}} \dots \dots (2)$$
Proof:

(1) Let $\varphi \in C_0^{\infty}(\mathbb{R}^d)$ Such that $\varphi = 1$ in the neighborhood of the ball of center 0 and radius r_1 . Let also $\tilde{\varphi} \in C_0^{\infty}(\mathbb{R}^d)$ Such that $\tilde{\varphi} \equiv 1$ in the neighborhood of φ , then we have

$$S_q u = \tilde{\varphi}(2^{-q}D)S_q u.$$

It is clear that

$$S_q u = \mathcal{F}^{-1} \big(\tilde{\varphi}(2^{-q}D) \mathcal{F}(S_q u) \big) = \mathcal{F}^{-1} (\tilde{\varphi}(2^{-q}D)) * S_q u.$$

Using the Fourier transform with a simple calculation, we get

$$\mathcal{F}^{-1}(\tilde{\varphi}(2^{-q}D)) = \int \tilde{\varphi}(2^{-q}\xi)e^{ix\xi}d\xi = 2^{qd}\int \tilde{\varphi}(\xi)e^{ix2^{q}\xi}d\xi$$
$$= 2^{qd}\mathcal{F}^{-1}(\tilde{\varphi}(\xi)) \coloneqq 2^{qd}h(2^{q}x).$$

This gives that

$$S_q u = 2^{qd} h(2^q \cdot) * S_q u.$$

Therefore

$$\partial^{\alpha} S_q u = 2^{q(d+|\alpha|)} \partial^{\alpha} h(2^q \cdot) * S_q u \dots \dots (3)$$

Taking the L^b norm of (3) and applying Young inequality for convolution (Lemma 1.1), we find with $\left(\frac{1}{b}+1=\frac{1}{c}+\frac{1}{a}\right)$, that $\left\|\partial^{\alpha}S_{q}u\right\|_{L^{b}} \leq 2^{q(d+|\alpha|)} \|\partial^{\alpha}h(2^{q}\cdot)\|_{L^{c}} \|S_{q}u\|_{L^{a}}$, $\leq 2^{q(d+|\alpha|)}2^{-q\frac{d}{c}} \|\partial^{\alpha}h\|_{L^{c}} \|S_{q}u\|_{L^{a}}$ $\leq 2^{q\left(|\alpha|+d\left(1-\frac{1}{c}\right)\right)} \|\partial^{\alpha}h\|_{L^{c}} \|S_{q}u\|_{L^{a}}$

$$\leq 2^{q\left(|\alpha|+d\left(\frac{z}{a}-\frac{z}{b}\right)\right)} \|\partial^{\alpha}h\|_{L^{c}} \|S_{q}u\|_{L^{a}}.$$

86

Therefore

$$sup_{|\alpha|=k} \left\| \partial^{\alpha} S_{q} u \right\|_{L^{b}} \leq 2^{q\left(k+d\left(\frac{1}{a}-\frac{1}{b}\right)\right)} \left\| \partial^{\alpha} h \right\|_{L^{c}} \left\| S_{q} u \right\|_{L^{a}} \dots \dots (4)$$

It is enough to prove that $\|\partial^{\alpha}h\|_{L^{c}} \leq C^{k}$. For this purpose, we use Lemma 1.2, then we have

$$\|\partial^{\alpha}h\|_{L^{c}} \leq \|\partial^{\alpha}h\|_{L^{1}} + \|\partial^{\alpha}h\|_{L^{\infty}} \dots \dots (5)$$

Now since $h = \mathcal{F}^{-1}\varphi$ and $\varphi \in \mathcal{C}_0^{\infty}(\mathbb{R}^d) \hookrightarrow S$, where S is the space of Schwartz, then we have $h \in S$, this gives by using Lemma 1.2, that h is bounded and $(1 + |\cdot|^2)^d \partial^{\alpha} h$ is also bounded.

 $|\cdot|^2)^d \partial^{\alpha} h \Big\|_{L^{\infty}}$

$$\leq C \left\| (1+|\cdot|^2)^d \partial^{\alpha} h \right\|_{L^{\infty}} \dots \dots (6)$$

Also

$$\begin{aligned} \|\partial^{\alpha}h\|_{L^{\infty}} &= \sup_{x} |\partial^{\alpha}h(x)| \leq \sup_{x} (1+|\cdot|^{2})^{d} |\partial^{\alpha}h| \\ &\leq C \left\| (1+|\cdot|^{2})^{d} \partial^{\alpha}h \right\|_{L^{\infty}} \dots \dots (7) \end{aligned}$$

Putting together (6) and (7) in (5), we get

$$\|\partial^{\alpha}h\|_{L^{c}} \leq C^{2} \|(1+|\cdot|^{2})^{d}\partial^{\alpha}h\|_{L^{\infty}} \leq C^{k}, \qquad k \in \mathbb{N}.$$

This gives in (4), that

$$\sup_{|\alpha|=k} \left\| \partial^{\alpha} S_{q} u \right\|_{L^{b}} \leq C^{k} 2^{q\left(k+d\left(\frac{1}{a}-\frac{1}{b}\right)\right)} \left\| S_{q} u \right\|_{L^{a}}.$$

This proves (1) of Lemma 2.1.

87

Proof of (2) of Lemma 2.1:

Let $\tilde{\varphi} \in \mathcal{C}_0^{\infty}(\mathbb{R}^d)$ Such that $\tilde{\varphi} = 1$ in the neighborhood of φ . Then we have

$$\dot{\Delta}_q u = \tilde{\varphi}(2^{-q}D)\dot{\Delta}_q u.$$

It is clear that

$$\mathcal{F}(\dot{\Delta}_q u) = \tilde{\varphi}(2^{-q}\xi)\mathcal{F}(\dot{\Delta}_q u),$$

and thus

$$\dot{\Delta}_{q}u(x) = \mathcal{F}^{-1}\left(\tilde{\varphi}(2^{-q}\xi)\mathcal{F}(\dot{\Delta}_{q}u)(\xi)\right)\dots\dots(8)$$

Where,

$$\tilde{\varphi}(2^{-q}\xi) = \sum_{|\alpha|=k} (i\xi)^{\alpha} |\xi|^{-2k} (-i\xi)^{\alpha} \tilde{\varphi}(2^{-q}\xi)$$

Putting this last inequality in (8), we get

$$\begin{split} \dot{\Delta}_{q}u(x) &= \sum_{|\alpha|=k} \mathcal{F}^{-1} \left((i\xi)^{\alpha} |\xi|^{-2k} (-i\xi)^{\alpha} \tilde{\varphi}(2^{-q}\xi) \mathcal{F} (\dot{\Delta}_{q}u)(\xi) \right) \\ &= \sum_{|\alpha|=k} \mathcal{F}^{-1} \left((i\xi)^{\alpha} |\xi|^{-2k} \tilde{\varphi}(2^{-q}\xi) \mathcal{F} (\partial^{\alpha} \dot{\Delta}_{q}u)(\xi) \right) \\ &= \sum_{|\alpha|=k} \mathcal{F}^{-1} \left((i\xi)^{\alpha} |\xi|^{-2k} \tilde{\varphi}(2^{-q}\xi) \right) * \partial^{\alpha} \dot{\Delta}_{q}u(x), \end{split}$$

where,

$$\begin{split} \mathcal{F}^{-1}\Big((i\xi)^{\alpha}|\xi|^{-2k}\tilde{\varphi}(2^{-q}\xi)\Big) &= \int \frac{(i\xi)^{\alpha}}{|\xi|^{2k}}\tilde{\varphi}(2^{-q}\xi)e^{ix\xi}d\xi\\ &= \int \frac{(i2^{q}\xi)^{\alpha}}{|2^{q}\xi|^{2k}}\tilde{\varphi}(\xi)e^{ix2^{qd}\xi}d\xi\\ &= 2^{q(d+|\alpha|-2k)}\int \frac{(i\xi)^{\alpha}}{|\xi|^{2k}}\tilde{\varphi}(\xi)e^{ix2^{qd}\xi}d\xi = 2^{q(d+|\alpha|-2k)}h_{k}(2^{q}x), \end{split}$$

88

where

$$h_k(2^q x) = \int \frac{(i\xi)^{\alpha}}{|\xi|^{2k}} \tilde{\varphi}(\xi) e^{ix2^{qd}\xi} d\xi.$$

Then

$$\dot{\Delta}_q u(x) = 2^{q(d+|\alpha|-2k)} h_k(2^q \cdot) * \partial^{\alpha} \dot{\Delta}_q u.$$

This gives by Young inequality for convolution, that

$$\left\|\dot{\Delta}_{q}u\right\|_{L^{a}} \leq 2^{q(d+|\alpha|-2k)} \|h_{k}(2^{q} \cdot)\|_{L^{1}} \left\|\partial^{\alpha}\dot{\Delta}_{q}u\right\|_{L^{a}}.$$

We have

$$\|h_k(2^q \cdot)\|_{L^1} = \int |h_k(2^q x)| dx = \int |h_k(y)| 2^{-qd} dy = 2^{-qd} \|h_k\|_{L^1}.$$

We recall that, $h = \mathcal{F}^{-1}\varphi$ and $\varphi \in \mathcal{C}_0^{\infty}(\mathbb{R}^d) \hookrightarrow S$, where S is the space of Schwartz, then we have $h \in S$, this gives by using Lemma 1.2, that *h* is bounded and $(1 + |\cdot|^2)^d \partial^{\alpha} h$ is also bounded. Therefore

 $\|h\|_{L^1} \leq C^k.$

Thus

$$\left\|\dot{\Delta}_{q}u\right\|_{L^{a}} \leq 2^{q(|\alpha|-2k)}C^{k}\left\|\partial^{\alpha}\dot{\Delta}_{q}u\right\|_{L^{a}}.$$

Thus

$$\sup_{|\alpha|=k} \left\| \partial^{\dot{\alpha}} \Delta_q u \right\|_{L^a} \ge C^{-k} 2^{qk} \left\| \dot{\Delta}_q u \right\|_{L^a}.$$

This is the desired result.

3. Conclusion

Bernstein inequality has been proved by J-Y-Chemin for any tempered distribution u. In this paper, we proved the inequality for a bloc dyadic $\dot{\Delta}_q u$ and $S_q u$. Our results show a strong support for the effect of mathematics and physical applications, for example, in the non-linear Navier-Stokes and Euler Boussinesq equations and the quasi-geostrophic equation.

89

References:

- [1] J.-Y.Chemin. (1998). Perfect incompressible fluids. Oxford University Press.
- [2] J.-Y.Chemin. (1999). Théorème d'unicité pour le système de Navier-Stokes tridimensionnel. J. Anal. Math., 77 no. 5, 27-50.
- [3] H. Bahouri, J-Y Chemin, and R. Danchin. (2011). Fourier analysis and nonlinear partial differential equations, volume 343. Springer Science & Business Media.
- [4] Q. Chen, C. Miao and Z. Zhang. (2007). A new Bernstein's inequality and the 2D dissipative quasi-geostrophic equation. Comm. Math. Phys., 271. No. 3, 821-838,.
- [5] *R Danchin.* (2012). Local theory in critical spaces for compressible viscous and Heat conductive gases.

Communications in Partial Differential Equations, 26:7-8, 1183-1233.

[6] P. G. Lemarié (2002). Recent developments in the Navier-Stokes problem, CRC Press.